

# DL Frequency Large Amplitude Modulation

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**DLS**  

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# Abstract

Intensity and frequency modulation by excitation current is well developed area of DL physics [1] with respect to rapid optical communication. We'll consider other subject of modulation related to TDLS operation.

The goal of present paper is to analyzed influence of DL excitation current modulation on recorded molecular line shape trace molecules detection.

Different modulation techniques were used in TDLS. However, these approaches were not optimal taking into account DL physical properties.

Resent TDLS operation regime includes DL frequency large amplitude modulation. Optimal modulation approach to achieve fundamental limit of minimum detectable absorption will be considered both based on theoretical and experimental investigations.

Achievement of minimum detectable absorption below  $10^{-7}$  will be demonstrated.

References:

[1] See for example: G.Morthier, P.Vankwikelberge, Handbook of Distributed Feedback Laser Diodes, Artech House, Inc., Boston, 1997, and references

# DL frequency modulation by current

DL frequency tuning by excitation current modulation is determined by two processes

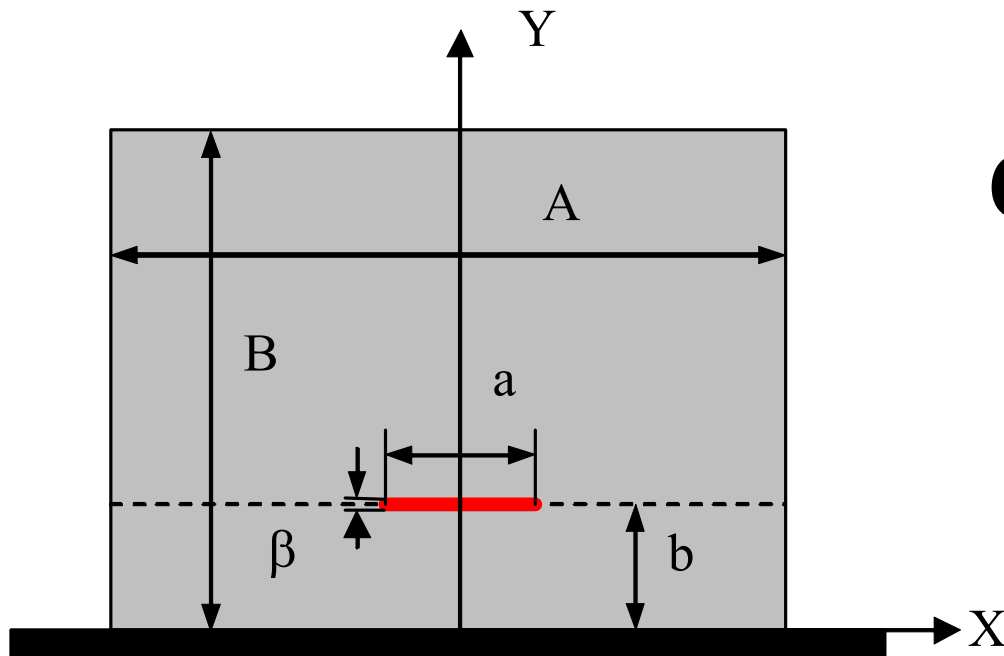
$$\frac{d\nu}{dI} = \left[ \frac{\partial \nu}{\partial T} \right]_I \frac{\partial T}{\partial I} + \left[ \frac{\partial \nu}{\partial I} \right]_T$$

Phonons system  
(temperature)

Electrons+photons  
system

# Temperature diffusion in DL

Cross-section of diode laser



Heat equation

$$C \frac{\partial \Delta T(x, y, t)}{\partial t} - \eta \left\{ \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right\} \Delta T(x, y, t) = W(x, y, t)$$

General solution heat equation

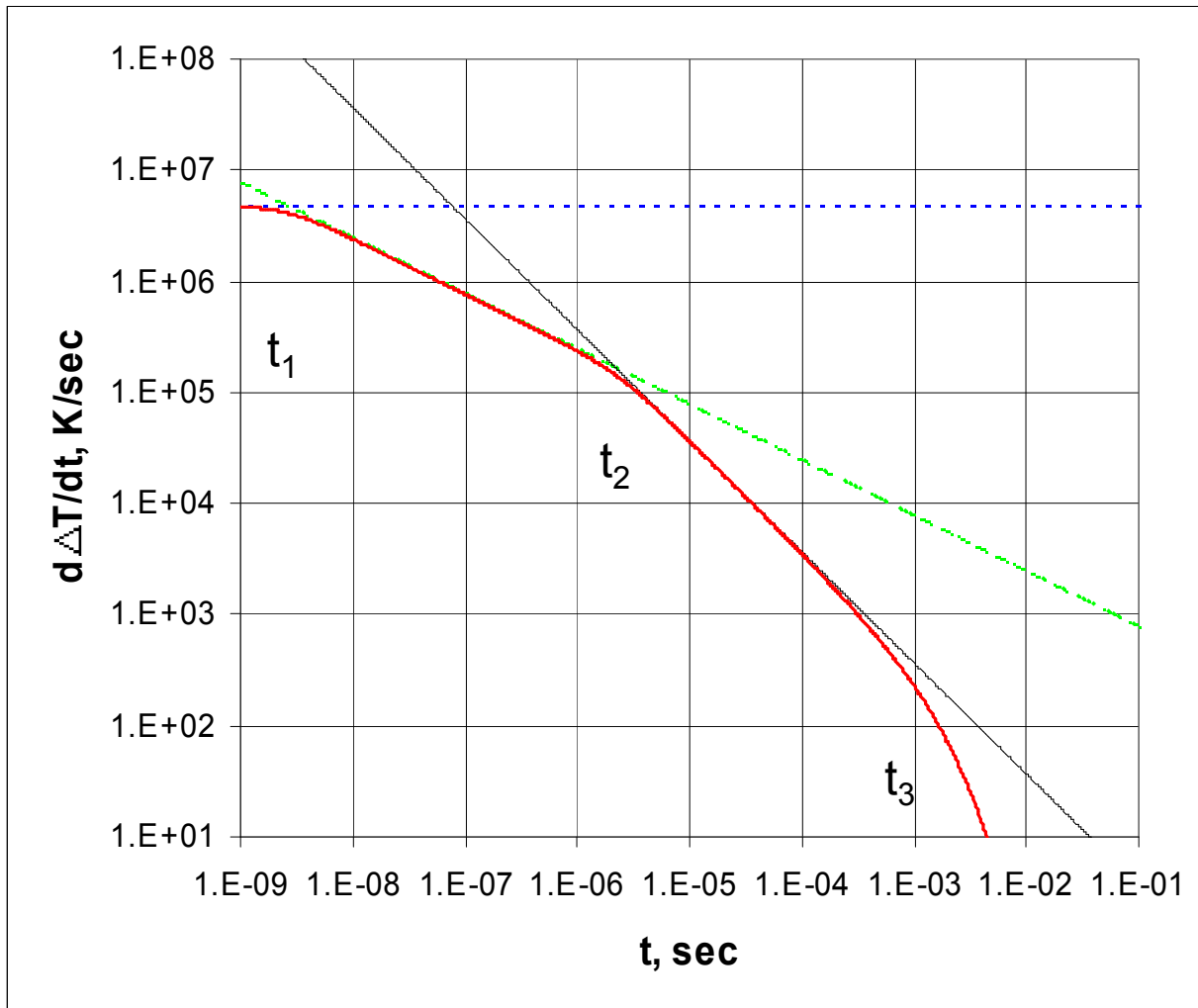
$$G(\vec{r}, \vec{\xi}, t) = \frac{1}{\sqrt{[4\pi Dt]^3}} \exp \left\{ -\frac{(\vec{r} - \vec{\xi})^2}{4Dt} \right\}$$

Boundary conditions

$$\Delta T \Big|_{y=0} = 0, \quad \frac{\partial \Delta T}{\partial y} \Big|_{y=B} = \frac{\partial \Delta T}{\partial x} \Big|_{x=\pm A/2} = 0$$

# Active area heating

Significant difference of the laser structure dimensions provides possibility to four different steps of heating process depending on thermal diffusion length  $L$ .



$L \ll \beta$  - homogenous heating of the DL active area with constant speed.  
 $\beta \ll L \ll a$  - thermal diffusion can be considered as one-dimensional process.  
 $a \ll L \ll B$  - two-dimensional process will take place.  
Finally, steady-state temperature distribution will be achieved after thermal diffusion length exceeds DL chip dimension -  $B$ .  
This leads to three characteristic times:  $t_1$ ,  $t_2$ ,  $t_3$

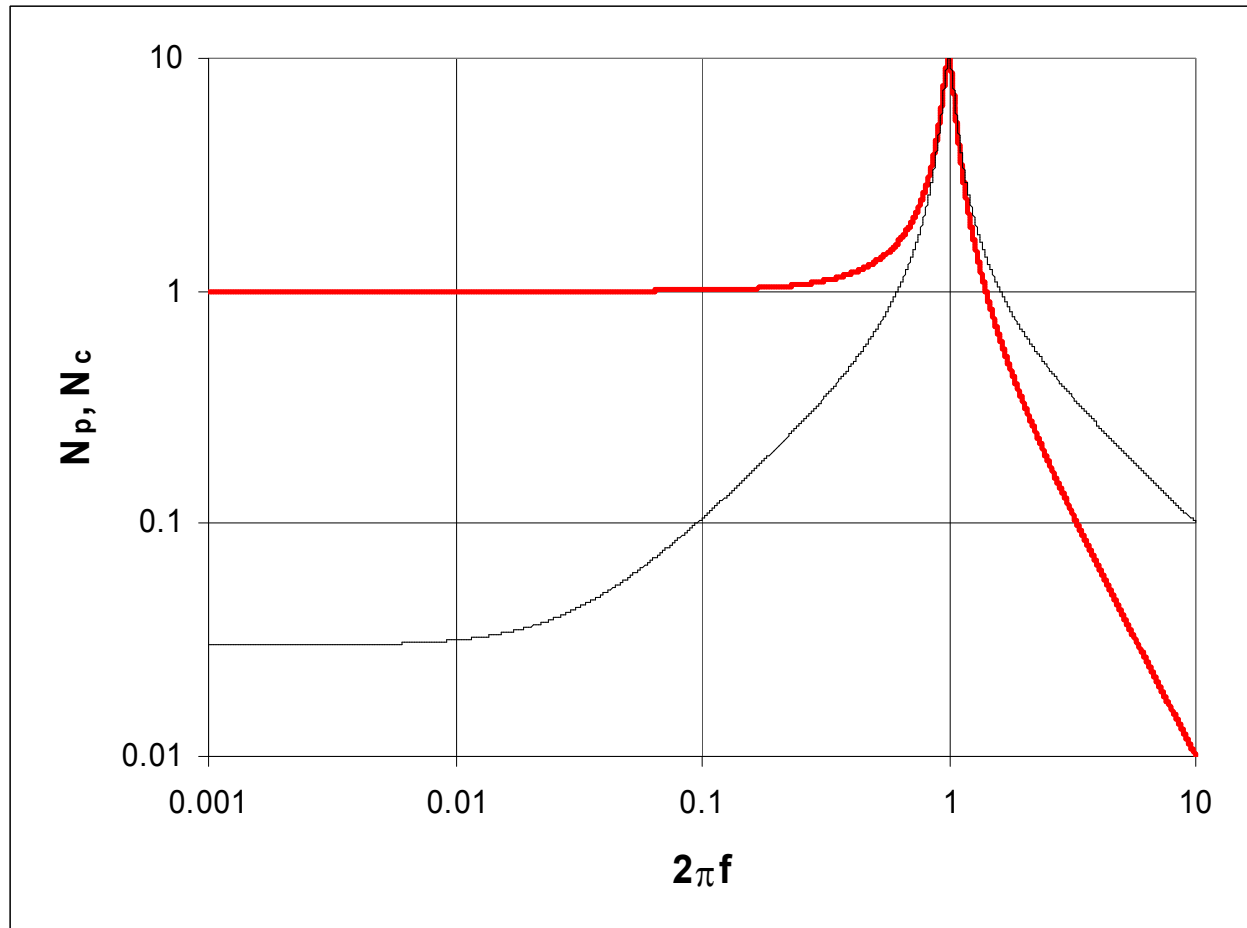
# Rate equations

Electrons + photons system. The behaviour of the electron subsystem in a diode laser is described by  $N_c$  - the number of non-equilibrium current carriers in the DL active area. Electromagnetic field is characterized by the number of photons  $N_p$  of m-th axial mode within which generation occurs.

$$\frac{dN_c}{dt} = \frac{I}{e} - g(N_c - N_G)N_p - \frac{N_c}{\tau_c}$$
$$\frac{dN_p}{dt} = g(N_c - N_G)N_p + \frac{CN_c}{\tau_s} - \frac{N_p}{\tau_p}$$

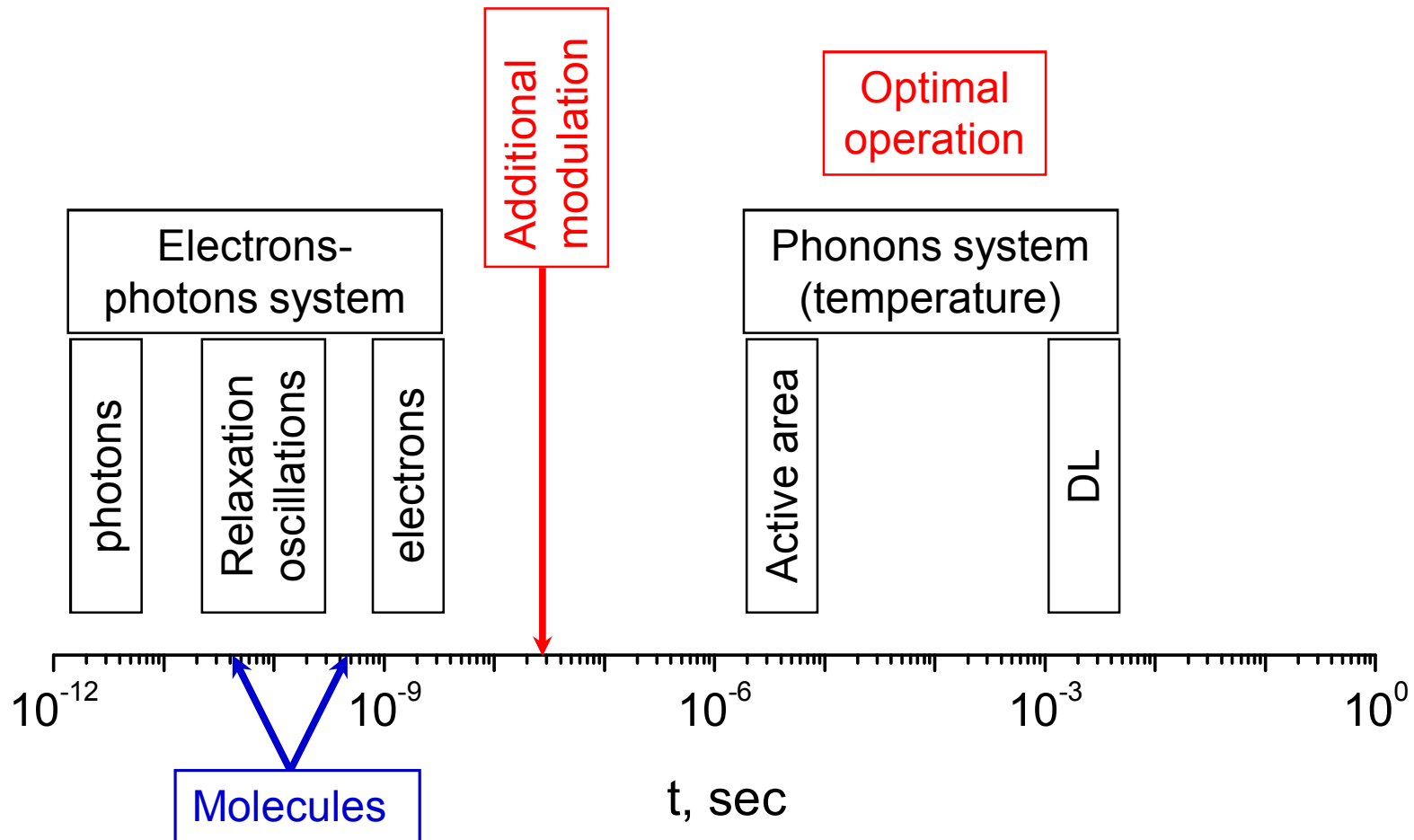
Here  $I$  is the excitation current;  $g$  - the constant proportional to the rate of the stimulated transition;  $N_G$  - the minimal number of the current carriers required for absorption compensation in a laser;  $\tau_c$  - the life time of a excess current carriers ( $\sim$  nsec);  $C$  - the contribution of a spontaneous emission into a given axial regime of a DL;  $\tau_s$  - spontaneous emission time;  $\tau_p$  - the life time of a photon in a resonator ( $\sim$  psec).

# Intensity and frequency modulation



DL power ( $N_p$ , black line) and frequency ( $N_c$ , red line) as function of modulation frequency. For a given case, frequency  $2\pi f$  was normalized to value  $\sqrt{gN_p/\tau_p}$  characterizing resonance known as “relaxation oscillations”.

# Physical processes in DL and optimal TDLS operation regime



In DL there are two systems: electrons+photons (fast) and phonons (slow)



# Frequency modulation

$$\vec{E}(t) = \vec{E}_0 \exp\left\{-i \int_{-\infty}^t 2\pi [c\nu + \Omega \cos(2\pi F\theta)] d\theta\right\} =$$

$$= \vec{E}_0 \exp\left\{-i \left[ 2\pi c\nu t + \frac{\Omega}{F} \sin(2\pi Ft) \right]\right\}$$

$c$  – light velocity;  
 $\nu = 1/\lambda$  [cm<sup>-1</sup>] - wavenumber;  
 $F$  [Hz] – modulation current frequency;  
 $\Omega$  [Hz] – light frequency deviation due to current modulation

$$\vec{E}_{slow}(t) = \vec{E}_0 \exp\left\{\frac{\Omega}{F} \exp(-i2\pi Ft)\right\} = \vec{E}_0 \sum_{n=0}^{\infty} \frac{1}{n!} \left[ \frac{\Omega}{F} \exp(-i2\pi Ft) \right]^n$$

$\Omega/F \ll 1$  – small fast modulation with small phase modulation

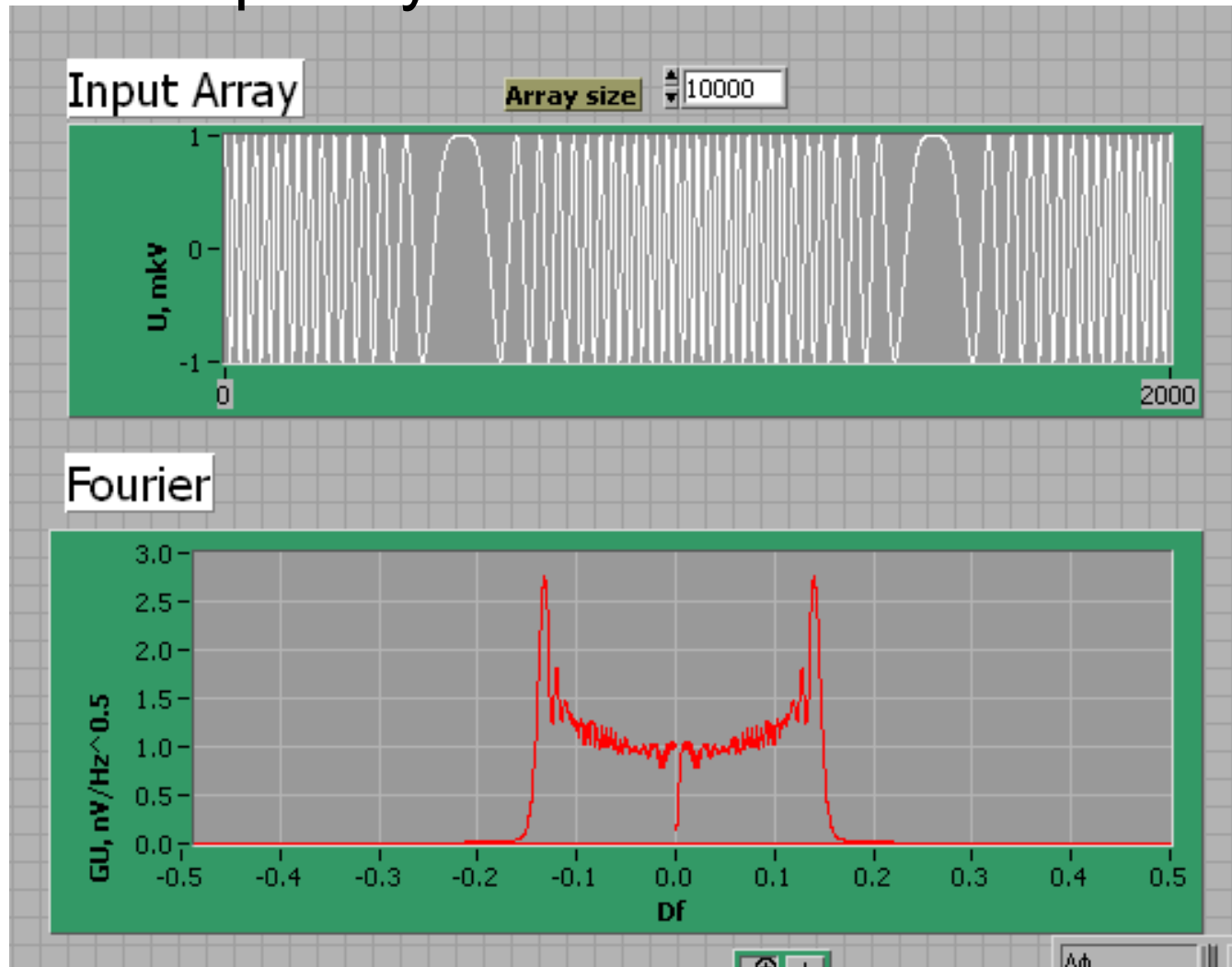
$\Omega/F \ll 1$  – large slow modulation with large phase modulation

Poisson distribution



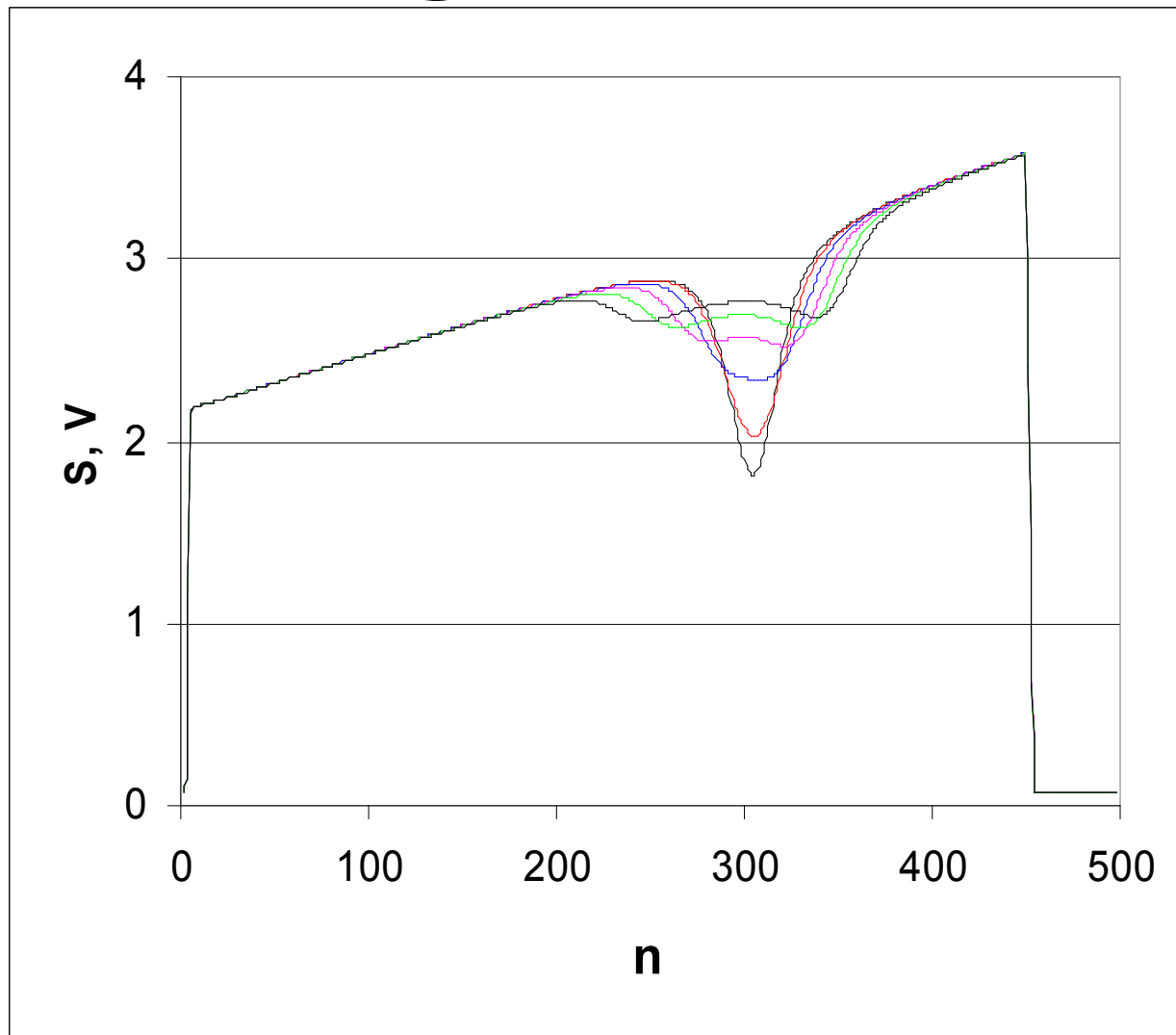
# Large amplitude frequency modulation

Frequency deviation  $\Omega \gg F$  - modulation frequency



Time dependence of slow field component determined by modulation and related spectrum demonstrating two peaks with similar amplitude

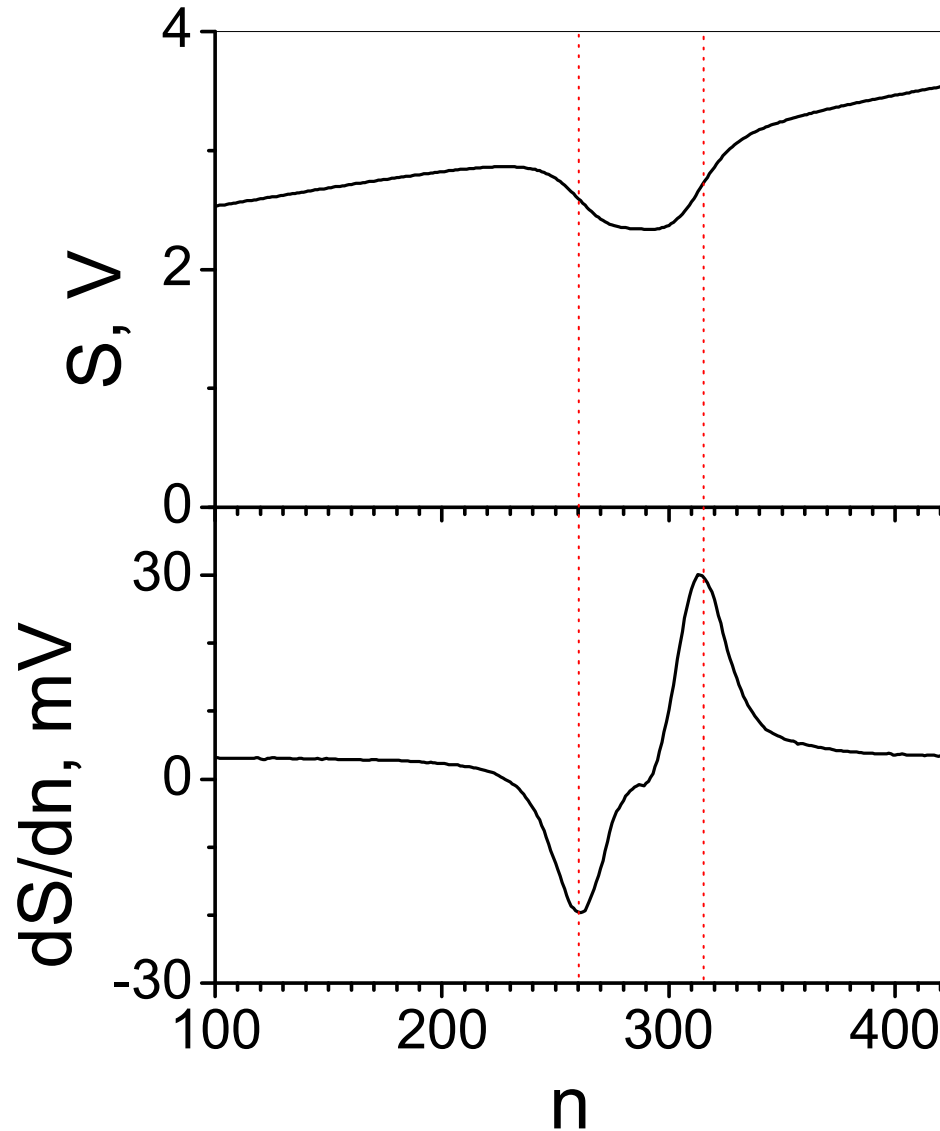
# Signal with modulation



Recorded signal with Doppler line shape of water vapor for several modulation amplitudes: 0, 4.7, 9.4, 14.1, 18.8, 23.5 mA;  $F = 25$  MHz

With modulation line broadening was observed with presence of two peaks having different amplitudes

# Line broadening due to modulation



Upper figure: signal with water Doppler line

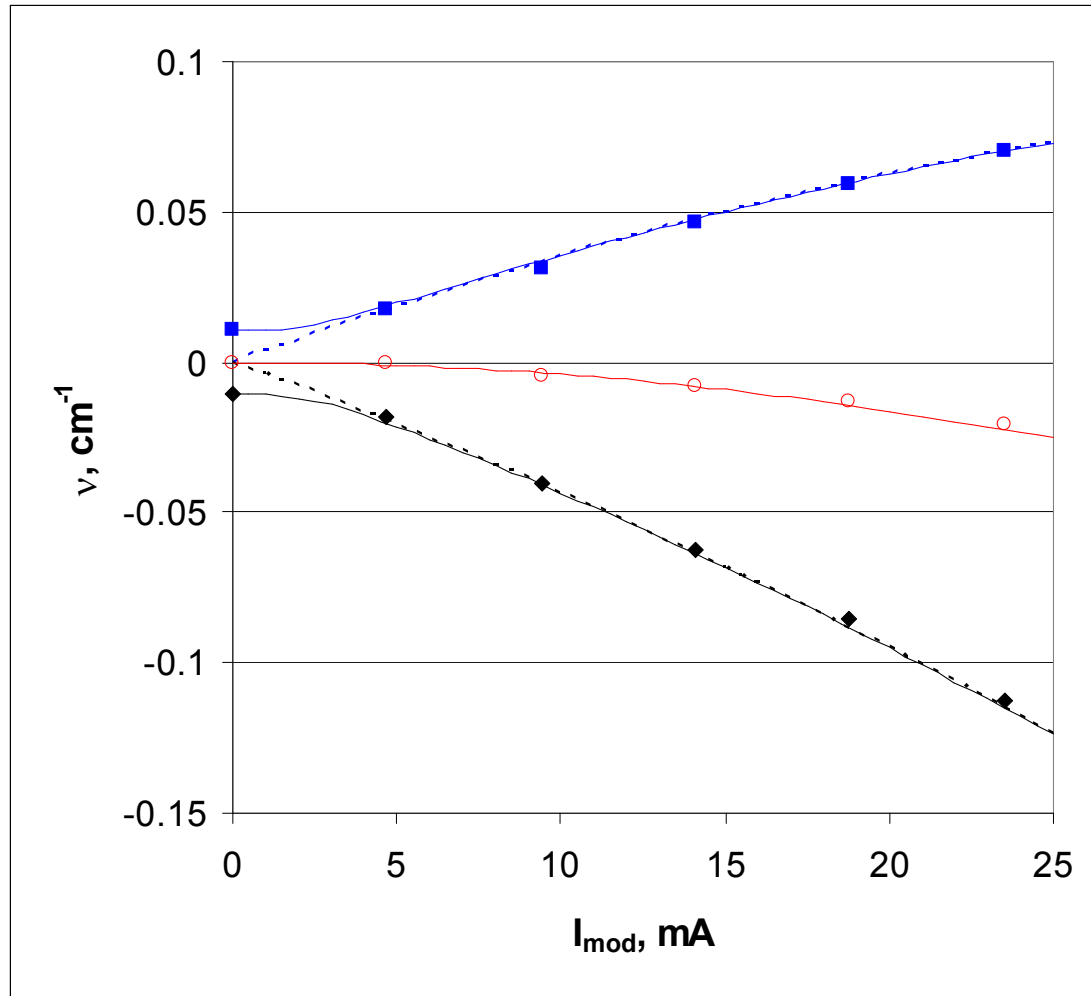
Lower figure: above spectrum derivative

Modulation frequency and amplitude were  $F = 25$  MHz and  $\Delta I = 9.4$  mA, respectively

Peaks on lower picture were used to determined line width

Line broadening asymmetry

# Line broadening due to modulation



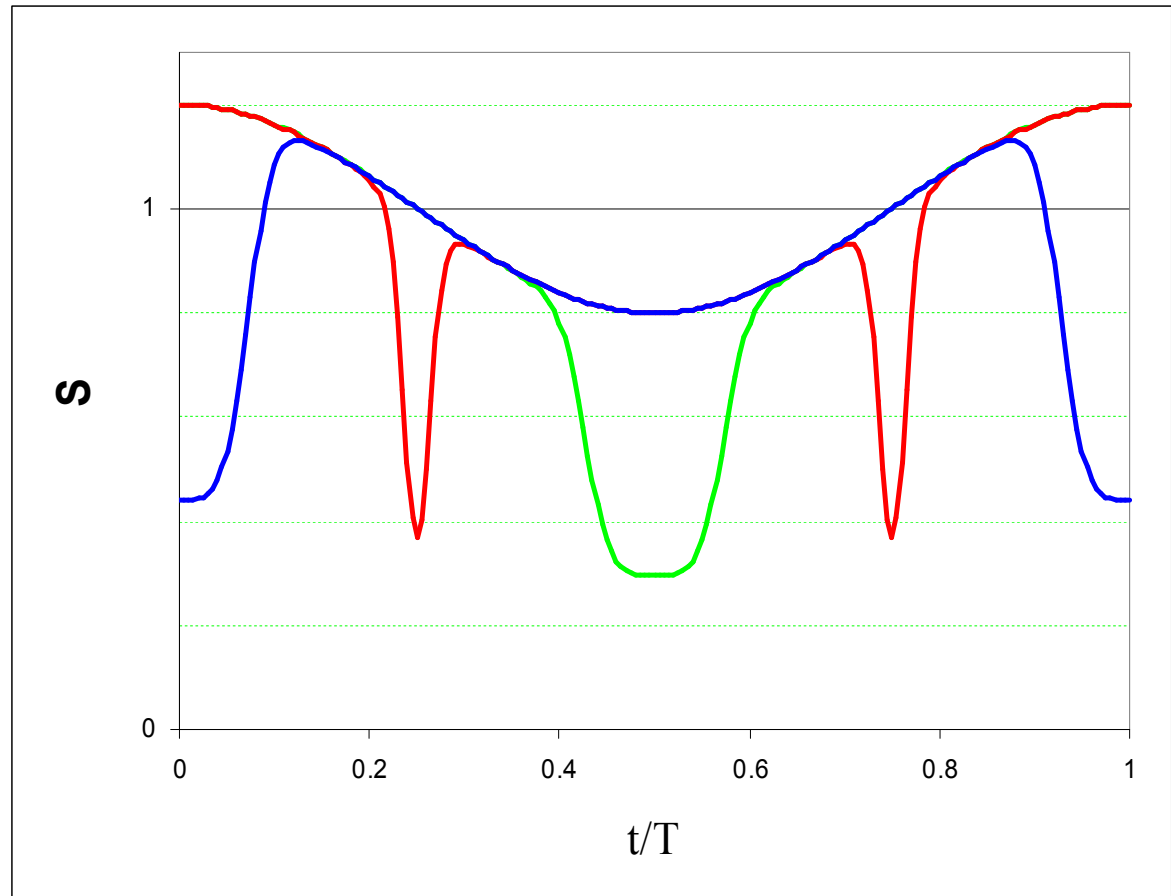
Broadening linear with  $\Delta I$  – fast electrons

Shift quadratic with  $\Delta I$  – slow temperature

# Spectral line with large modulation

Electrons+photons system of DL is fast enough in comparison with modulation frequency -  $F$ . Hence, moment DL wavenumber -  $\nu(t)$  can be introduced. Molecules are also fast with respect with  $F$  and there absorption can be determined as function of  $\nu(t)$ .

Fig. presents examples of transmitted light time dependence when line center corresponds to max, min, and center of  $\nu(t)$ .



Detection system is slow in comparison with  $F$  and record mean value of signals presented on Fig.

# Spectral line shapes for large frequency modulation

